

# DFG-Schwerpunktprogramm 1324

„Extraktion quantifizierbarer Information aus komplexen Systemen“

## An iterative reconstruction scheme for photoacoustic imaging

Y. Dong, T. Görner, S. Kunis

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# An iterative reconstruction scheme for photoacoustic imaging

Yiqiu Dong<sup>†</sup>      Torsten Görner<sup>‡</sup>      Stefan Kunis<sup>§</sup>

Recovery of image data from photoacoustic measurements asks for the inversion of the spherical mean value operator. In contrast to direct inversion methods for specific geometries, we consider a semismooth Newton scheme to solve a total variation regularized least squares problem. During the iteration, each matrix vector multiplication is realized in an efficient way using a recently proposed spectral discretization of the spherical mean value operator. All theoretical results are illustrated by numerical experiments.

*Key words and phrases* : spherical mean operator, fast Fourier transform, total variation regularization, photoacoustic imaging.

*2010 AMS Mathematics Subject Classification* : 65T50, 44A12, 92C55.

## 1 Introduction

Analogously to the Radon transform for computerized tomography, the spherical mean value operator is the crucial ingredient in photoacoustic imaging [31, 21, 4]. Recovery of image data from photoacoustic measurements hence asks for the inversion of this operator and this problem has been studied recently in [9, 3, 1, 21, 28] and references therein. For specific geometries, direct reconstruction algorithms are discussed in [14, 25, 24, 13, 2, 8, 23, 11] and in contrast to point-like detectors in the above model, variants for integrating detectors have been studied in [12, 5, 26, 34, 15, 33].

In this paper we consider the case of center points located at an arbitrary submanifold where direct reconstruction formulae are unknown. In order to reconstruct images from photoacoustic measurements stably, regularization techniques are utilized. In this direction and with the aim of preserving significant edges in images, we consider the total variation regularization, which was proposed in [29] and is well known in image restoration. Then, we set up an iterative method to solve the total variation regularized least squares problem based on the Fenchel-duality and inexact semismooth Newton techniques following the approach in

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[17]. In each iteration, we apply the recently proposed algorithm [10] for the fast and accurate computation of spherical means.

The structure of our paper is as follows. Section 2 reviews the Cauchy problem for the wave equation and its relation to the spherical mean value operator. After discretization of this operator via trigonometric polynomials, we discuss the algorithmic properties of this discrete spherical mean value operator. The following section sets up the total variation regularized least squares problem ( $P_0$ ) and studies a tight relaxation for which the semismooth Newton method [27, 16] is applied. Several numerical experiments are reported in Section 4 and we conclude our findings in the last section.

## 2 Photoacoustic imaging, spherical means, and discretization

In many topics of photoacoustic imaging we have to deal with the Cauchy problem for the wave equation

$$\begin{aligned} \partial_t^2 p(\mathbf{x}, t) - \nu_s^2(\mathbf{x}) \Delta p(\mathbf{x}, t) &= 0 \quad \text{for } (\mathbf{x}, t) \in \mathbb{R}^d \times (0, \infty), \\ p(\mathbf{x}, 0) &= f(\mathbf{x}) \quad \text{for } \mathbf{x} \in \mathbb{R}^d, \\ \partial_t p(\mathbf{x}, 0) &= 0 \quad \text{for } \mathbf{x} \in \mathbb{R}^d, \end{aligned} \tag{2.1}$$

where  $d \in \mathbb{N}$ ,  $d \geq 2$ , is the spatial dimension,  $p$  is the pressure and  $\nu_s$  is the speed of sound within the medium. In general, the speed of sound depends on the spatial variable  $\mathbf{x}$ , but many setups allow to assume homogeneity and consequently, after rescaling, we suppose  $\nu_s = 1$ . It is well known [7, §13, eq. 13-15] that for a sufficiently smooth function  $f : \mathbb{R}^d \rightarrow \mathbb{R}$  the solution of (2.1) can be given as

$$p(\mathbf{x}, t) = \begin{cases} \frac{1}{\Gamma(\frac{d}{2})2^{d-1}} (t^{-1} \frac{\partial}{\partial t})^{\frac{d}{2}} \int_0^t \frac{r}{\sqrt{t^2-r^2}} (r^{-1} \frac{\partial}{\partial r})^{\frac{d-2}{2}} r^{d-2} (\mathcal{M}f)(\mathbf{x}, r) dr & \text{for even } d, \\ \frac{\sqrt{\pi}}{\Gamma(\frac{d}{2})2^{\frac{d-1}{2}}} \frac{\partial}{\partial t} (t^{-1} \frac{\partial}{\partial t})^{\frac{d-3}{2}} (t^{d-2} (\mathcal{M}f)(\mathbf{x}, t)) & \text{for odd } d, \end{cases}$$

where the spherical mean value operator  $\mathcal{M}$  is defined by

$$\mathcal{M}f(\mathbf{x}, r) := \frac{1}{\omega_{d-1}} \int_{\mathbb{S}^{d-1}} f(\mathbf{x} + r\boldsymbol{\xi}) d\sigma(\boldsymbol{\xi}), \quad (\mathbf{x}, r) \in \mathbb{R}^d \times \mathbb{R}.$$

Here,  $\mathbb{S}^d := \{\mathbf{x} \in \mathbb{R}^d : |\mathbf{x}|_2^2 := \sum_{j=1}^d x_j^2 = 1\}$  denotes the  $d-1$  dimensional sphere,  $\sigma$  denotes the surface measure on  $\mathbb{S}^{d-1}$  and  $\omega_{d-1} := \sigma(\mathbb{S}^{d-1})$ . The goal is to recover  $f$ , if  $p(\mathbf{x}, t)$ ,  $(\mathbf{x}, t) \in \Omega \times \mathbb{R}$ , is known, where  $\Omega \subset \mathbb{R}^d$  is some submanifold surrounding the region of interest. In particular for spatial dimension  $d = 3$ , we have to find  $f(\mathbf{x})$ ,  $\mathbf{x} \in \mathbb{R}^3$ , from the data

$$p(\mathbf{x}, t) = \frac{\partial}{\partial t} (t (\mathcal{M}f)(\mathbf{x}, t)), \quad \mathbf{x} \in \Omega \subset \mathbb{R}^3, \quad t \in \mathbb{R}.$$

In this paper, we limit our considerations to recovering a function from its spherical means with iterative methods. For this purpose, the efficient and accurate computation of spherical means from given function values on some grid is essential. We follow our recent approach in [10] and restrict to functions supported on  $[-\frac{1}{2}, \frac{1}{2}]^d =: \mathbb{T}^d$ , which are approximated by trigonometric polynomials. The spherical mean value operator is bounded from  $L^p(\mathbb{T}^d)$  to

$L^p(\mathbb{T}^d \times [0, 1], d\mathbf{y}r^{d-1} dr)$ , and allows for a decomposition in eigenfunctions  $e_{\mathbf{k}} : \mathbb{T}^d \rightarrow \mathbb{C}$ ,  $e_{\mathbf{k}}(\mathbf{x}) = e^{2\pi i \mathbf{k} \mathbf{x}}$ ,  $\mathbf{k} \in \mathbb{Z}^d$ ,

$$\mathcal{M}e_{\mathbf{k}}(\mathbf{y}, r) = e_{\mathbf{k}}(\mathbf{y}) \frac{\Gamma\left(\frac{d}{2}\right) J_{\frac{d-2}{2}}(2\pi|\mathbf{k}|_2 r)}{(\pi|\mathbf{k}|_2 r)^{\frac{d-2}{2}}},$$

where  $J$  denotes the Bessel function of first kind. For some discretization parameter  $n \in \mathbb{N}$ , the function  $f : \mathbb{T}^d \rightarrow \mathbb{R}$  is typically given by discrete values  $f(\mathbf{x})$  on a regular grid  $\mathbf{x} \in X \subset \mathbb{T}^d$ ,

$$X := \left\{ \left( \frac{2j_1+1-n}{2n}, \dots, \frac{2j_d+1-n}{2n} \right)^\top, \mathbf{j} \in [0, n)^d \cap \mathbb{Z}^d \right\},$$

and the spherical means  $\mathcal{M}f(\mathbf{y}, r)$  have to be computed for scattered center points  $\mathbf{y} \in Y \subset \mathbb{T}^d$  and radii  $r \in R \subset [0, 1]$ . To this end, we compute discrete Fourier coefficients

$$\hat{f}_{\mathbf{k}} := \sum_{\mathbf{x} \in X} f(\mathbf{x}) e^{-2\pi i \mathbf{k} \mathbf{x}}, \quad \mathbf{k} \in J_n := \left[ -\frac{n}{2}, \frac{n}{2} \right)^d \cap \mathbb{Z}^d,$$

by one ordinary fast Fourier transform (FFT) in a first step. Subsequently, we compute for each radius  $r \in R$  auxiliary coefficients

$$\tilde{h}_{\mathbf{k}, r} := \hat{f}_{\mathbf{k}} \frac{\Gamma\left(\frac{d}{2}\right) J_{\frac{d-2}{2}}(2\pi|\mathbf{k}|_2 r)}{(\pi|\mathbf{k}|_2 r)^{\frac{d-2}{2}}}, \quad \mathbf{k} \in J_n,$$

and evaluate

$$\mathcal{M}f(\mathbf{y}, r) \approx g(\mathbf{y}, r) := \sum_{\mathbf{k} \in J_n} \tilde{h}_{\mathbf{k}, r} e^{2\pi i \mathbf{k} \mathbf{y}}, \quad \mathbf{y} \in Y,$$

by nonequispaced FFTs. In matrix vector notation, this sums up to

$$\mathbf{g} = \mathbf{M} \mathbf{f},$$

with a discrete spherical mean value operator  $\mathbf{M} : \mathbb{R}^N \rightarrow \mathbb{C}^M$ ,  $N = n^d$ ,  $M = M_1 M_2$ ,  $M_1 = |Y|$ ,  $M_2 = |R|$ ,

$$\mathbf{M} = \begin{pmatrix} \mathbf{A} \mathbf{D}_1 \mathbf{F} \\ \mathbf{A} \mathbf{D}_2 \mathbf{F} \\ \vdots \\ \mathbf{A} \mathbf{D}_{M_2} \mathbf{F} \end{pmatrix},$$

with the following nonequispaced Fourier matrix  $\mathbf{A} \in \mathbb{C}^{M_1 \times N}$ , diagonal matrices  $\mathbf{D}_j \in \mathbb{R}^{N \times N}$ ,  $j = 1, \dots, M_2$ , and a Fourier matrix  $\mathbf{F} \in \mathbb{C}^{N \times N}$ ,

$$\begin{aligned} (\mathbf{A})_{\mathbf{y}, \mathbf{k}} &= e^{2\pi i \mathbf{k} \mathbf{y}}, & \mathbf{y} \in Y, \mathbf{k} \in J_n, \\ (\mathbf{D}_j)_{\mathbf{k}, \mathbf{k}} &= \frac{\Gamma\left(\frac{d}{2}\right) J_{\frac{d-2}{2}}(2\pi\|\mathbf{k}\|_2 r_j)}{(\pi\|\mathbf{k}\|_2 r_j)^{\frac{d-2}{2}}}, & \mathbf{k} \in J_n, \\ (\mathbf{F})_{\mathbf{k}, \mathbf{x}} &= e^{-2\pi i \mathbf{k} \mathbf{x}}, & \mathbf{k} \in J_n, \mathbf{x} \in X. \end{aligned}$$

An immediate consequence is given by

**Lemma 2.1.** *With the above definitions, we have*

$$\mathbf{M}^* \mathbf{M} = \mathbf{F}^* \left( \sum_{j=1}^{M_2} \mathbf{D}_j \mathbf{T} \mathbf{D}_j \right) \mathbf{F}$$

with the hermitian multilevel Toeplitz matrix  $\mathbf{T} \in \mathbb{C}^{N \times N}$ ,

$$(\mathbf{T})_{\mathbf{k}, \mathbf{l}} = \sum_{\mathbf{y} \in Y} e^{2\pi i(\mathbf{k}-\mathbf{l})\mathbf{y}}.$$

*Proof.* We rewrite the above definition as

$$\mathbf{M} = \begin{pmatrix} \mathbf{A} & & \\ & \ddots & \\ & & \mathbf{A} \end{pmatrix} \begin{pmatrix} \mathbf{D}_1 \\ \vdots \\ \mathbf{D}_{M_2} \end{pmatrix} \mathbf{F}$$

and use

$$(\mathbf{A}^* \mathbf{A})_{\mathbf{k}, \mathbf{l}} = \sum_{\mathbf{y} \in Y} e^{2\pi i(\mathbf{k}-\mathbf{l})\mathbf{y}}$$

■

As detailed in [10], the above approach leads to an algorithm of complexity  $\mathcal{O}(M_2(N \log N + M_1))$  in general using the nonequispaced FFT [20]. For the two-dimensional case, i.e.  $n \times n$  images,  $M_1 = M_2 = \mathcal{O}(n)$  detectors and radii, this leads to a semi-fast method of complexity  $\mathcal{O}(n^3 \log n)$ . In case  $d = 3$ , i.e.  $n \times n \times n$  volumes,  $M_1 = \mathcal{O}(n^2)$  detectors and  $M_2 = \mathcal{O}(n)$  radii, this can be improved by the recently proposed butterfly sparse FFT [32, 22] and leads to an algorithm of complexity  $\mathcal{O}(n^3 \log n)$ .

We note in passing that the considered trigonometric interpolation is complex valued also for real valued functions, but can easily be made real valued by extending the discrete Fourier coefficients  $\hat{f}_{\mathbf{k}}$  to  $[-\frac{n}{2}, \frac{n}{2}]^d \cap \mathbb{Z}^d \supset J_n$  appropriately. Subsequently, we consider this real valued discrete spherical mean value operator and for ease of notation the two-dimensional case.

### 3 TV-Based Iterative Method

We now turn to the reconstruction problem

$$\mathbf{M} \mathbf{f} = \mathbf{g}, \tag{3.1}$$

where  $\mathbf{g} \in \mathbb{R}^M$  is the vector of discrete spherical mean values,  $\mathbf{f} \in \mathbb{R}^N$  is a real valued image obtained from a two-dimensional  $n$ -by- $n$  pixel-array on a regular grid by concatenation in the usual columnwise fashion with  $N = n^2$ , and  $\mathbf{M} \in \mathbb{R}^{M \times N}$  is the discrete spherical mean value operator. The problem of reconstructing the image  $\mathbf{f}$  from the measurements  $\mathbf{g}$  is known to be ill-posed. Hence, regularization techniques based on prior information on  $\mathbf{f}$  are utilized to get stable reconstruction processes. In this direction, total variation (TV) regularization [29] had great success, which is defined by

$$\|\mathbf{v}\|_{\text{TV}} := \sum_{k=1}^N \|\nabla \mathbf{v}\|_k = \sum_{k=1}^N \sqrt{|\nabla_x \mathbf{v}|_k^2 + |\nabla_y \mathbf{v}|_k^2}.$$



Here, we use the abbreviation  $[\mathbf{p}]_k = (p_k, p_{N+k})^\top$ ,  $\mathbf{p} \in \mathbb{R}^{2N}$ ,  $k = 1, \dots, N$ , and define the discrete gradient operator  $\nabla \in \mathbb{R}^{2N \times N}$  by  $[\nabla \mathbf{v}]_k := ((\nabla_x \mathbf{v})_k, (\nabla_y \mathbf{v})_k)^\top$  with the forward differences

$$(\nabla_x \mathbf{v})_k = \begin{cases} \mathbf{v}_{s+1,t} - \mathbf{v}_{s,t}, & \text{if } s < n, \\ 0, & \text{if } s = n, \end{cases} \quad (\nabla_y \mathbf{v})_k = \begin{cases} \mathbf{v}_{s,t+1} - \mathbf{v}_{s,t}, & \text{if } t < n, \\ 0, & \text{if } t = n, \end{cases}$$

and  $k = sn + t$  with  $(s, t) \in \{1, 2, \dots, n\} \times \{1, 2, \dots, n\}$  as an index in the regular pixel-array. Combining the TV regularization with an  $l^2$ -data-fitting term, we reconstruct the image  $\mathbf{f}$  by solving the minimization problem

$$\min_{\mathbf{f} \in \mathbb{R}^N} \mathcal{J}(\mathbf{f}), \quad \mathcal{J}(\mathbf{f}) := \frac{1}{2} \|\mathbf{M}\mathbf{f} - \mathbf{g}\|_2^2 + \alpha \|\mathbf{f}\|_{\text{TV}}, \quad (P_0)$$

where  $\alpha > 0$  is the regularization parameter to control the trade-off between a good fitness to  $\mathbf{g}$  and the smoothness from the TV term.

The TV term is nondifferentiable and this is responsible for preserving edges in images but also poses an algorithmic challenge. In order to overcome the difficulty, referring to [17] we utilize Fenchel duality and introduce associated inexact semismooth Newton techniques. Applying the Fenchel-Legendre duality calculus, we derive the Fenchel-dual of  $(P_0)$

$$\begin{aligned} \sup_{\mathbf{p} \in \mathbb{R}^{2N}} & -\frac{1}{2} \|\mathbf{M}^* \mathbf{g} - \text{div} \mathbf{p}\|_{\mathbf{M}}^2 + \frac{1}{2} \|\mathbf{g}\|_2^2, \\ \text{s.t.} & \quad \|\mathbf{p}\|_2 \leq \alpha, \quad \text{for all } k = 1, \dots, N, \end{aligned} \quad (P_0^*)$$

where  $\|\mathbf{v}\|_{\mathbf{M}}^2 := \langle (\mathbf{M}^* \mathbf{M})^{-1} \mathbf{v}, \mathbf{v} \rangle = \sum_{k=1}^N \mathbf{v}_k \cdot ((\mathbf{M}^* \mathbf{M})^{-1} \mathbf{v})_k$ , and the divergence operator is given by  $\text{div} = -\nabla^\top$ . Due to the nontrivial kernel of the divergence operator, the solution of  $(P_0^*)$  is not unique. To overcome the non-uniqueness of the solution of  $(P_0^*)$ , we propose a dual regularization

$$\begin{aligned} \sup_{\mathbf{p} \in \mathbb{R}^{2N}} & -\frac{1}{2} \|\mathbf{M}^* \mathbf{g} - \text{div} \mathbf{p}\|_{\mathbf{M}}^2 + \frac{1}{2} \|\mathbf{g}\|_2^2 - \frac{\gamma}{2\alpha} \sum_{k=1}^N \|\mathbf{p}\|_k^2, \\ \text{s.t.} & \quad \|\mathbf{p}\|_k \leq \alpha, \quad \text{for all } k = 1, \dots, N, \end{aligned} \quad (P^*)$$

where  $\gamma > 0$  is the dual regularization parameter. In order to understand the effect of the dual regularization on the original minimization problem, we apply the Fenchel-Legendre calculus once more and find that the dual of  $(P^*)$  is given by

$$\min_{\mathbf{f} \in \mathbb{R}^N} \frac{1}{2} \|\mathbf{M}\mathbf{f} - \mathbf{g}\|_2^2 + \sum_{k=1}^N (\Phi_\gamma(\nabla \mathbf{f}))_k, \quad (P)$$

where

$$(\Phi_\gamma(\nabla \mathbf{f}))_k := \begin{cases} \frac{\alpha}{2\gamma} \|[\nabla \mathbf{f}]_k\|_2^2, & \text{if } \|[\nabla \mathbf{f}]_k\|_2 < \gamma, \\ \alpha (\|[\nabla \mathbf{f}]_k\|_2 - \frac{\gamma}{2}), & \text{if } \|[\nabla \mathbf{f}]_k\|_2 \geq \gamma. \end{cases}$$

Here,  $\Phi_\gamma$  is a Huber function [19], which smooths locally the TV term in order to obtain the differentiability of  $(P)$  and the uniqueness of the dual solution  $\mathbf{p}$ .

Based on the first order optimality condition, the solution  $(\bar{\mathbf{f}}, \bar{\mathbf{p}})$  satisfies

$$\mathbf{M}^* \mathbf{M} \bar{\mathbf{f}} + \text{div} \bar{\mathbf{p}} = \mathbf{M}^* \mathbf{g}, \quad (3.2)$$

$$\max\{\gamma, |[\nabla \bar{\mathbf{f}}]_k|_2\}[\bar{\mathbf{p}}]_k = -\alpha[\nabla \bar{\mathbf{f}}]_k, \quad \text{for } k = 1, \dots, N.$$

Due to the presence of the max-operator, the system is not smooth but semismooth. We are able to solve it by the semismooth Newton technique, see [27, 16]. The semismooth Newton step is given by

$$\mathbf{f}^{l+1} = \mathbf{f}^l + \delta, \quad \mathbf{H}^l \delta = \mathbf{w}^l \quad (3.3)$$

with

$$\begin{aligned} \mathbf{H}^l &= \mathbf{M}^* \mathbf{M} + \nabla^\top \mathbf{D}(\mathbf{m}_{\gamma^l})^{-1} \left[ \alpha \mathbf{I}_N - \chi_{\mathcal{A}^l} \mathbf{D}(\mathbf{p}^l) \mathbf{N}(\nabla \mathbf{f}^l) \right] \nabla, \\ \mathbf{w}^l &= -\mathbf{M}^*(\mathbf{M} \mathbf{f}^l - \mathbf{g}) - \alpha \nabla^\top \mathbf{D}(\mathbf{m}_{\gamma^l})^{-1} \nabla \mathbf{f}^l. \end{aligned}$$

Here,  $\mathbf{I}_N \in \mathbb{R}^{N \times N}$  is the identity matrix,  $\mathbf{D}(\mathbf{m}_{\gamma^l}) \in \mathbb{R}^{2N \times 2N}$  is a diagonal matrix with the vector  $\mathbf{m}_{\gamma^l} \in \mathbb{R}^{2N}$ ,  $(\mathbf{m}_{\gamma^l})_k = (\mathbf{m}_{\gamma^l})_{N+k} = \max\{\gamma, |[\nabla \mathbf{f}^l]_k|_2\}$ , as the main diagonal. Moreover, we define the set  $\mathcal{A}^l = \{k : |[\nabla \mathbf{f}^l]_k|_2 > \gamma\}$  and the thresholding operator as diagonal matrix  $\chi_{\mathcal{A}} \in \mathbb{R}^{2N \times 2N}$ ,

$$(\chi_{\mathcal{A}})_{k,k} = (\chi_{\mathcal{A}})_{N+k,N+k} = \begin{cases} 1, & \text{if } k \in \mathcal{A}, \\ 0, & \text{if } k \in \mathcal{A}^c. \end{cases}$$

In addition, with  $\mathbf{v} = (\mathbf{v}_1^\top, \mathbf{v}_2^\top)^\top \in \mathbb{R}^{2N}$ , we set

$$\mathbf{N}(\mathbf{v}) = \begin{bmatrix} \mathbf{D}(\mathbf{v}_1) & \mathbf{D}(\mathbf{v}_2) \\ \mathbf{D}(\mathbf{v}_1) & \mathbf{D}(\mathbf{v}_2) \end{bmatrix} \in \mathbb{R}^{2N \times 2N}.$$

In general, the matrix  $\mathbf{H}^l$  is not symmetric, but at a solution it is symmetric positive semi-definite. In order to avoid problems due to these potential deficiencies during the iterations, we utilize the modifications as introduced in [17], whenever these are necessary. In addition, since there is no guarantee that  $\mathbf{H}^l$  is positive definite, a solution  $\delta$  in (3.3) may not exist or may not be unique in case of existence. As a remedy, we add a small multiple of the identity matrix to the system matrix and then solve it by biconjugate gradient stabilized (BICGSTAB) algorithm [30]. Similar as in [17] the whole algorithm can be shown to converge at a superlinear rate provided that  $\mathbf{f}^0$  is sufficiently close to the solution.

## 4 Numerical experiments

In this section, we provide numerical results to study the behavior of our method with respect to its image reconstruction capability and its computational efficiency. In our numerics, when solving the minimization problem (3.2), we set  $\gamma = 10^{-3}$ , and we stop the Newton iteration as soon as the initial residual is reduced by a factor of  $10^{-4}$ . In each Newton step, the BICGSTAB iteration is stopped as soon as the relative norm of the residual of the primal-dual system (3.2) at  $(\mathbf{f}^l, \mathbf{p}^l)$  drops below

$$\begin{aligned} \text{tol}^{l+1} &:= 10^{-3} \cdot \min \left\{ \left( \frac{\|\mathbf{r}^l\|_2}{\|\mathbf{r}^0\|_2} \right)^{\frac{3}{2}}, \frac{\|\mathbf{r}^l\|_2}{\|\mathbf{r}^0\|_2} \right\}, \\ \mathbf{r}^l &:= \left( \begin{array}{c} \mathbf{M}^* \mathbf{M} \mathbf{f}^l + \text{div} \mathbf{p}^l - \mathbf{M}^* \mathbf{g} \\ (\max\{\gamma, |[\nabla \mathbf{f}^l]_k|_2\}[\mathbf{p}^l]_k + \alpha[\nabla \mathbf{f}^l]_k)_{k=1, \dots, N} \end{array} \right). \end{aligned}$$

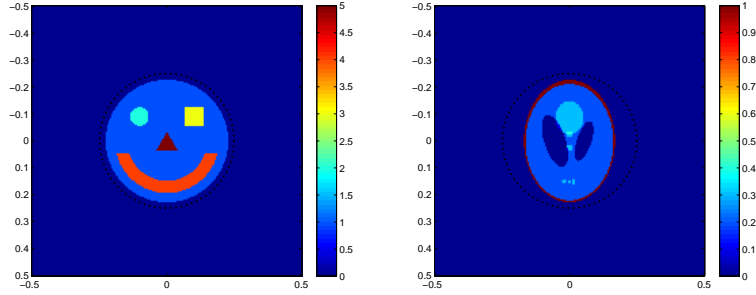


Figure 4.1: The phantoms Smiley and Shepp Logan.

Table 4.1: For each iteration  $l$  we show the residual  $\|\mathbf{r}^l\|_2$  of (3.2), the objective function value  $\mathcal{J}(\mathbf{f})$  of  $(P_0)$  and the number of the BICGSTAB iterations  $l_{\text{inner}}$ .

$l$	Smiley Face			Shepp Logan Phantom		
	$\ \mathbf{r}^l\ _2$	$\mathcal{J}(\mathbf{f}^l)$	$l_{\text{inner}}$	$\ \mathbf{r}^l\ _2$	$\mathcal{J}(\mathbf{f}^l)$	$l_{\text{inner}}$
0	0.0960	0.0923	-	0.0212	0.0084	-
1	0.0016	0.0081	30	9.66e-4	0.0028	24
2	7.53e-4	0.0068	21	4.79e-4	0.0024	19
3	7.04e-4	0.0066	16	1.50e-5	0.0024	43
4	4.88e-4	0.0065	13	1.68e-6	0.0023	44
5	1.07e-4	0.0065	16			
6	1.96e-7	0.0065	45			

Furthermore, for the comparison of computational efficiency, all simulations are run in MATLAB 7.11.0 (R2010b) on a laptop equipped with a P8700 2.53GHz CPU and 4GByte main memory.

*Example 1.* In this example, we give the reconstructed results from the spherical mean values of the phantoms shown in Figure 4.1 by solving the TV regularization model  $(P_0)$ . Here, the phantoms are assumed to be supported in the disk with the radius less than 0.25, and there are  $M_1 = 80$  detectors (shown as black dots in Figure 4.1) uniformly distributed on the surrounding circle. Since in this case the spherical mean values vanish for  $r > 0.5$ , the given data  $\mathbf{g}$  include the spherical mean values with  $0 \leq r \leq 0.5$ , which are discretized as  $M_2 = 100$  linearly equally spaced points. Hence, in this example we have  $M = 8000$ . In addition, the image resolution of the phantoms are 100-by-100, i.e.  $N = n^2 = 10^4$ . In Figure 4.2, we show the spherical mean values  $\mathbf{g}$  and the least squares solutions of (3.1) by solving  $\min_{\mathbf{f}} \frac{1}{2} \|\mathbf{M}\mathbf{f} - \mathbf{g}\|_2^2$ , which is utilized as the initial value in our primal-dual method for solving the TV-model.

When solving the minimization problem  $(P_0)$ , the regularization parameter is set to  $\alpha = 10^{-5}$ , and the stopping rule of the Newton iteration are reached after  $l = 6$  and  $l = 4$  iterations with the CPU time 222.1s and 173.3s, respectively. Figure 4.3 shows the reconstructed images in a few iterations. Comparing the final results with the original phantoms in Figure 4.1 and

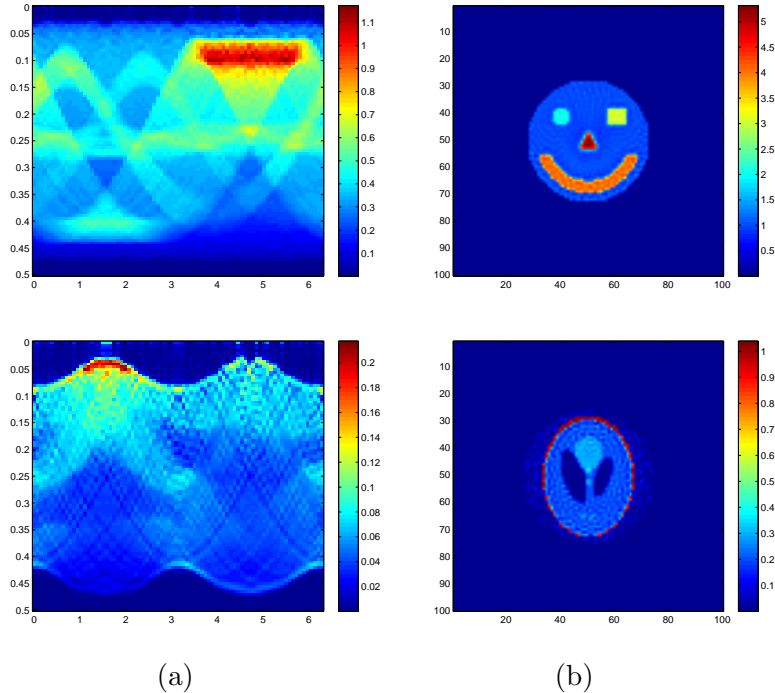


Figure 4.2: (a) The data  $\mathbf{g} \in \mathbb{R}^{M_1 \times M_2}$ , (b) the least squares solutions of (3.1), row 1: for the smiley face; row 2: for the Shepp Logan phantom.

the least squares solutions, we observe that the method based on the TV-model performs very good reconstructions. Furthermore, to illustrate the convergence behavior of our method for solving (3.2), in Table 4.1 we list the residuals of the system (3.2), the objective function values of  $(P_0)$ , and the number of the BICGSTAB algorithm in each Newton step. The decrease of the residual implies a locally superlinear convergence of the iterations.

*Example 2.* In order to show the computational efficiency, we compare our primal-dual method with an alternating minimization algorithm which utilizes the splitting technique to solve the TV-model  $(P_0)$ , see [18]. In the proposed algorithm, a quadratic penalty term is introduced in  $(P_0)$ , i.e.,

$$\min_{\mathbf{f} \in \mathbb{R}^N} \frac{1}{2} \|\mathbf{M}\mathbf{f} - \mathbf{g}\|_2^2 + \frac{\beta}{2} \|\mathbf{f} - \mathbf{u}\|_2^2 + \alpha \|\mathbf{u}\|_{\text{TV}}.$$

Then, the reconstructed image  $\mathbf{f}$  is obtained by solving the following two minimization problems

$$\begin{aligned} \mathbf{f}^l &= \operatorname{argmin}_{\mathbf{f} \in \mathbb{R}^N} \frac{1}{2} \|\mathbf{M}\mathbf{f} - \mathbf{g}\|_2^2 + \frac{\beta}{2} \|\mathbf{f} - \mathbf{u}^{l-1}\|_2^2, \\ \mathbf{u}^l &= \operatorname{argmin}_{\mathbf{u} \in \mathbb{R}^N} \frac{\beta}{2} \|\mathbf{f}^l - \mathbf{u}\|_2^2 + \alpha \|\mathbf{u}\|_{\text{TV}}, \end{aligned}$$

alternately. The objective function of the first optimization problem is quadratic, and it is solved by the conjugate gradient method. The second minimization problem is working on the image domain, and the Fenchel-duality-based Chambolle's method [6] is used to solve it

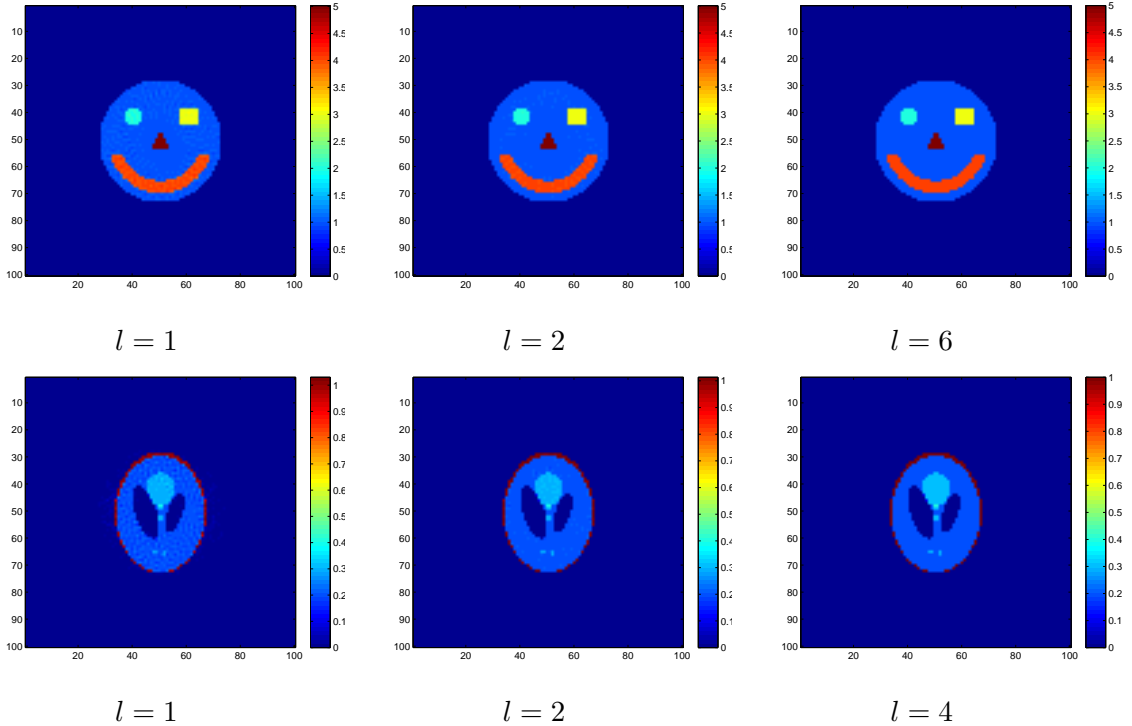


Figure 4.3: The reconstructed result in the Newton iteration during solving the TV-model.

as suggested in [18] with  $\beta = 3 \cdot 10^{-3}$ . Since both of the methods are solving approximations of  $(P_0)$ , to compare the computational efficiency, in Figure 4.4 we show the plots of the objective function values of  $(P_0)$  versus iterations with the same  $\alpha$ . In this example, we stop the splitting method after 200 iterations and it spends 6648.6 seconds. But the final function value is still larger than that from our method after 6 Newton iterations and 222.1s CPU time. It is easy to see that our method is much more efficient. In addition, we also give the reconstructed images obtained by both methods. Since the alternating algorithm and our primal-dual method both aim at minimizing the TV-model  $(P_0)$ , they obtain similar results. But comparing the intensity range of the results, our method performs more exact reconstruction, and the shrinkage in the result by the alternating minimization algorithm is due to the quadratic penalty function technique.

*Example 3.* Now, we test our algorithm for reconstructing the data corrupted by white Gaussian noise with different noise levels. In each case, we choose a reasonable parameter  $\alpha$  for it. In Figure 4.5, we give the data  $\mathbf{g}$  corrupted by 5% and 10% white Gaussian noise, respectively, and the corresponding reconstructed images by solving  $\min_{\mathbf{f}} \|\mathbf{M}\mathbf{f} - \mathbf{g}\|_2^2$ . The results obtained by our algorithm from solving the TV-model  $(P_0)$  are shown in Figure 4.6. Since the parameter  $\alpha$  controls the trade-off between a good fit of  $\mathbf{g}$  and a smoothness requirement due to the total variation regularization, to show the influence of the parameter  $\alpha$  on the reconstructions, we list two results with different values of  $\alpha$  for each noise level in Figure 4.6. It is seen that with larger  $\alpha$ , the reconstructed images are smoother, and the intensity range of the images are reduced obviously, which can be seen from the color bar on the right side of the images. However, with smaller  $\alpha$  the results include more details, but at

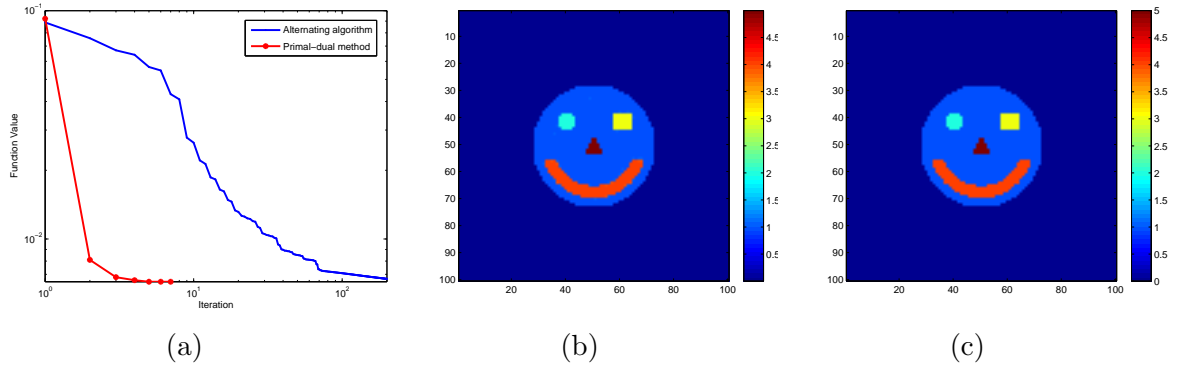


Figure 4.4: (a) The objective function values of  $(P_0)$  versus iterations, (b) the result from the alternating minimization algorithm proposed in [18], (c) the result from our primal-dual method.

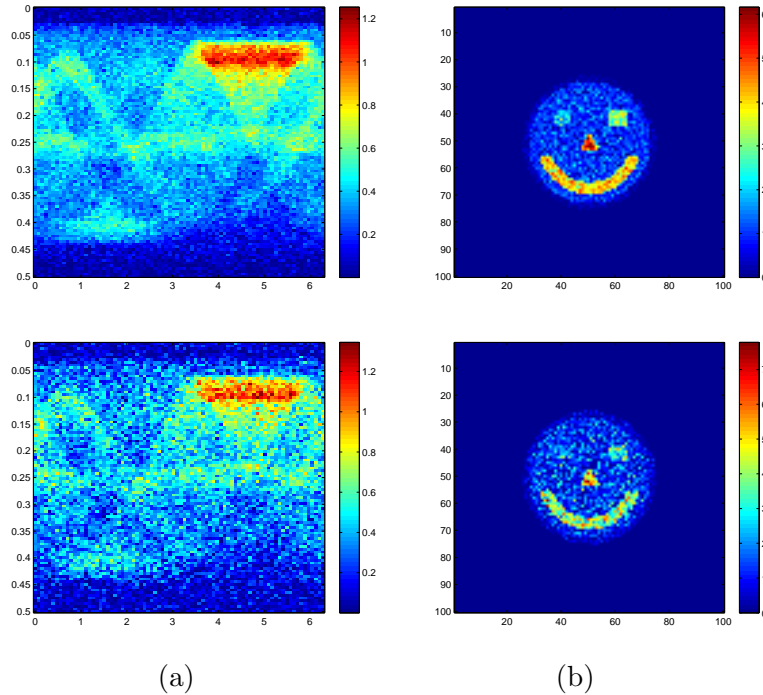


Figure 4.5: (a) With the noise level 5%; (b) with the noise level 10%. Row 1: the data  $\mathbf{g}$  with white Gaussian noise; row 2: the least squares solutions of (3.1).

the same time some noise is left.

In addition, in Table 4.2 we list the number of the Newton iterations, the total number of the BICGSTAB iterations and the CPU time. Here, the least squares solutions cannot supply good estimate as the initial vector  $\mathbf{f}^0$ , see Figure 4.5. Although in all our tests the numbers of Newton steps reach the maximum of 10 iterations, the reconstructed results are already acceptable. Based on our numerical experiments, allowing more iterations yields no

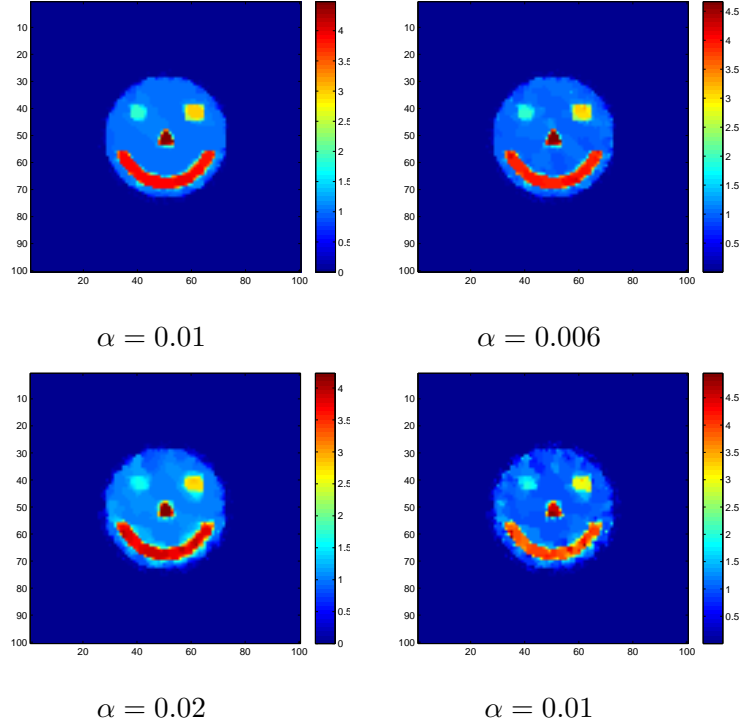


Figure 4.6: The reconstructed results by solving the TV-model with different parameter values and noise levels. Row 1: for noise level 5%; row 2: for noise level 10%.

Table 4.2: For different noise levels and  $\alpha$ , the number of Newton iterations  $l^{\text{all}}$ , the total number of the BICGSTAB iterations  $l_{\text{inner}}^{\text{all}}$ , and the CPU time  $t$ .

noise	5%		10%	
	0.01	0.006	0.02	0.01
$l^{\text{all}}$	10	10	10	10
$l_{\text{inner}}^{\text{all}}$	124	120	124	168
$t$	179.6	170.4	189.8	242.9

significant effect on the results.

*Example 4.* Since in many practical cases the detectors are only able to be located on a limited region in this example we consider that the detectors are uniformly distributed on the upper half circle instead of whole circle. In Figure 4.7, we give the data  $\mathbf{g}$ , and the least squares solution of (3.1), which is set as the initial guess of  $\mathbf{f}$  when solving the TV-model. In addition, Figure 4.7 also shows the reconstructed image with  $\alpha = 10^{-4}$  by our algorithm based on the TV-model and the plot of the residuals of the primal-dual system equations (3.2). We note that in this limited view case solving the TV-model by our primal-dual method still gives a good reconstruction.

*Example 5.* In this example, we consider the image size of phantoms as 200-by-200, i.e.

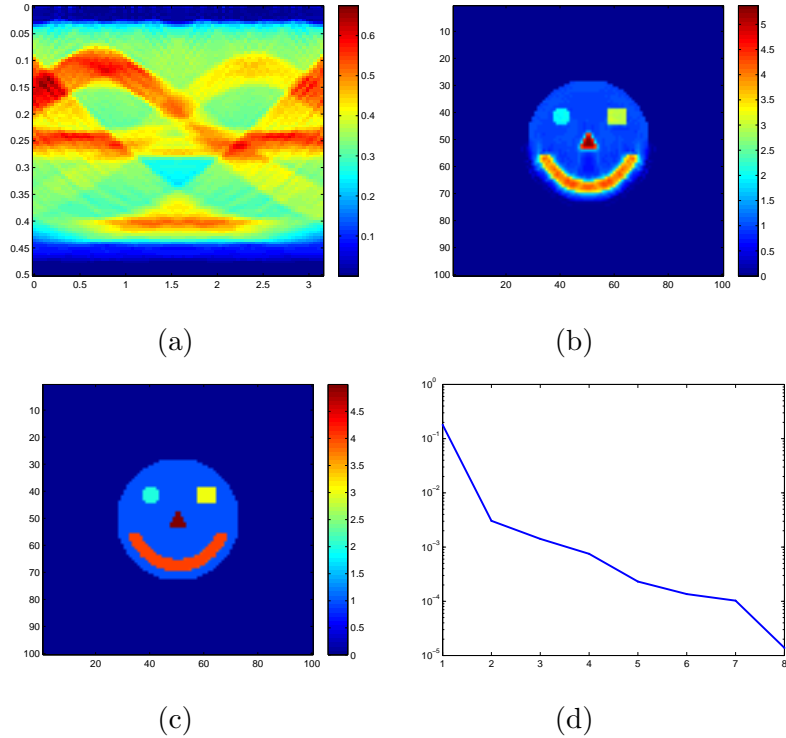


Figure 4.7: (a) The data  $\mathbf{g}$ , (b) the reconstructed result by solving  $\min \|\mathbf{M}\mathbf{f} - \mathbf{g}\|_2^2$ , (c) the result by solving  $(P_0)$ , (d) the residuals of the system equations (3.2) versus iterations.

$n = 200$ , and there are still  $M_1 = 80$  detectors uniformly distributed on the unit circle. But the range of the parameter  $r$  in the spherical mean values is discretized as  $M_2 = 200$  linearly equally spaced points. In Figure 4.8, we give the data  $\mathbf{g}$ , the least squares results, and the reconstructed images with  $\alpha = 6 \cdot 10^{-5}$  by solving the TV-model. By our method, the stopping rule is reached after  $l = 5$  and  $l = 10$  Newton iterations with the CPU time 2337.9s and 2431.5s, respectively.

## 5 Summary

We introduced a novel iterative reconstruction method in photoacoustic imaging which does not rely on a specific geometry of the detectors. Based on a dedicated discretization of the spherical mean value operator and a total variation regularizer, our method performs reasonable efficiently and preserves important features like edges in the reconstruction.

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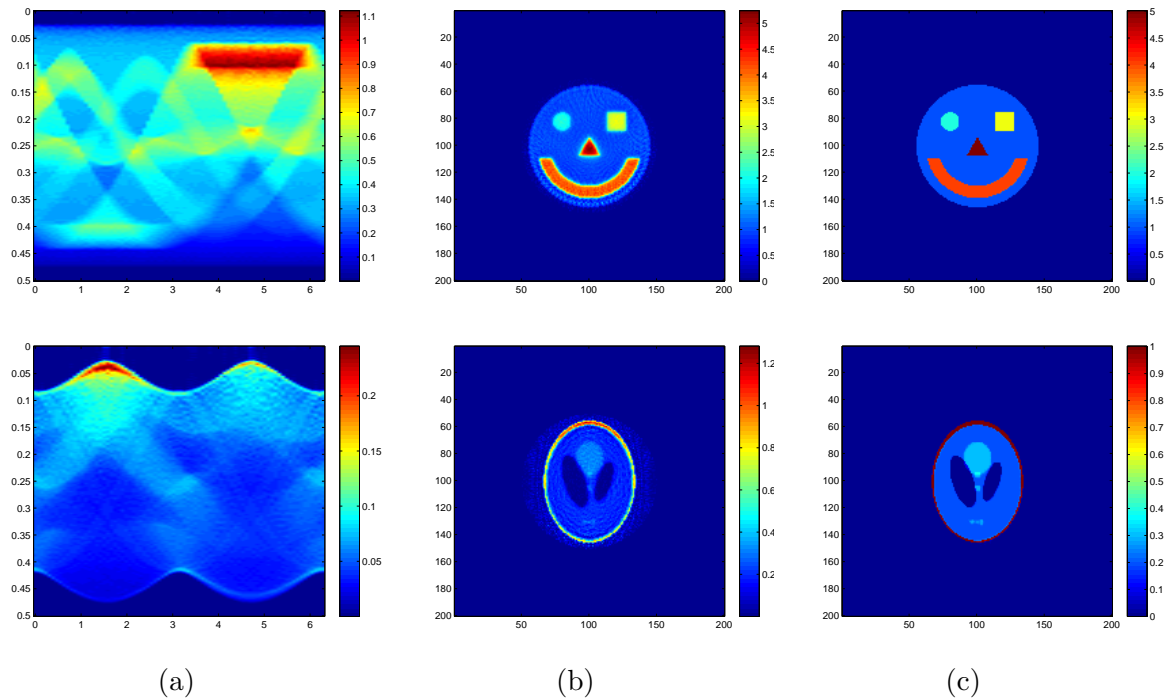


Figure 4.8: (a) The data  $g$ , (b) the least squares results, (c) the reconstructed result by solving  $(P_0)$ .

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